

# Entanglement and secret-key-agreement capacities of bipartite quantum interactions and read-only memory devices

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# Bipartite quantum interactions

Bipartite unitary interactions are the most elementary many-body interactions. Due to unavoidable interaction with environment, study of bipartite noisy interactions is pertinent.

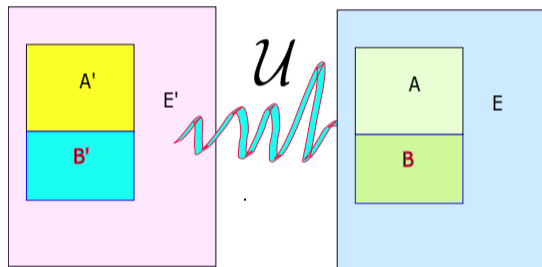


Figure: Systems of interest  $A'$  and  $B'$  interacting in presence of the bath  $E'$ .

- $\mathcal{U}$  is unitary transformation corresponding to underlying interaction Hamiltonian  $\hat{H}$  among  $A', B', E'$ .
- Before action of interaction Hamiltonian  $\hat{H}$ :  $\omega_{A'B'} \otimes \tau_{E'}$ , where bath  $E'$  is in some fixed state and uncorrelated to  $A'B'$ .
- After action of  $\hat{H}$ :

$$\rho_{ABE} := \mathcal{U}_{A'B'E' \rightarrow ABE}(\omega_{A'B'} \otimes \tau_{E'}).$$

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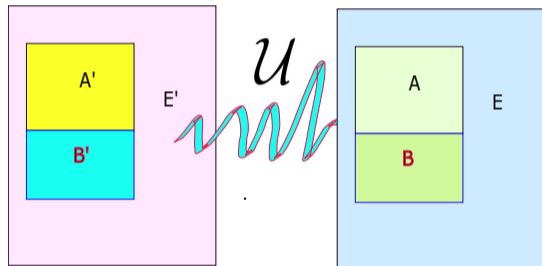


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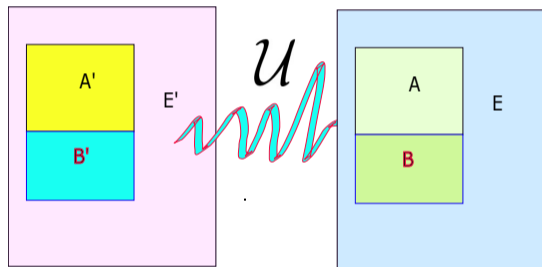


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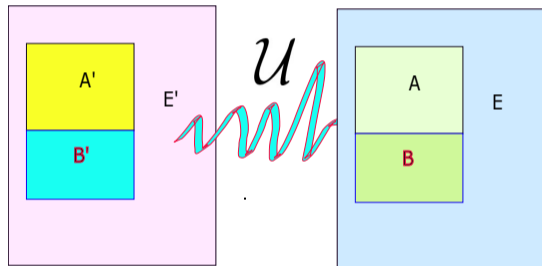


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## Bidirectional quantum channel

A bipartite quantum channel  $\mathcal{N}_{A'B' \rightarrow AB}$  is a completely positive, trace-preserving map that transforms composite system  $A'B'$  to  $AB$ .

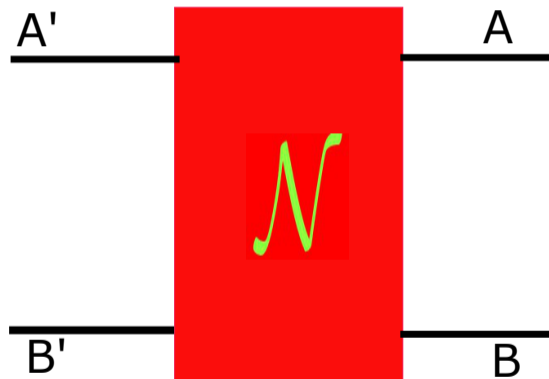


Figure: Two parties of interest: Alice holds  $A'$ ,  $A$  and Bob holds  $B'$ ,  $B$ .

- When  $A'$ ,  $A$  are held by Alice and  $B'$ ,  $B$  are held by Bob, bipartite channel  $\mathcal{N}$  is called bidirectional channel.
- It corresponds to noisy bipartite interaction, when bath is inaccessible.
- For all input state  $\omega_{A'B'}$ :

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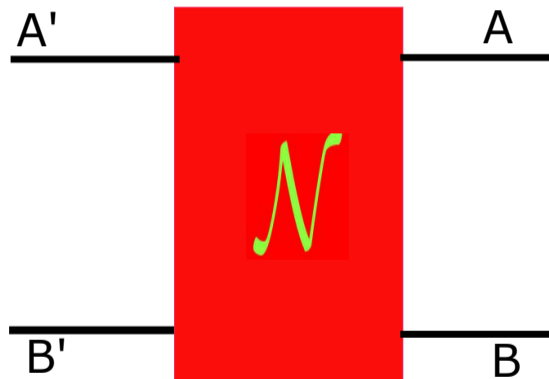


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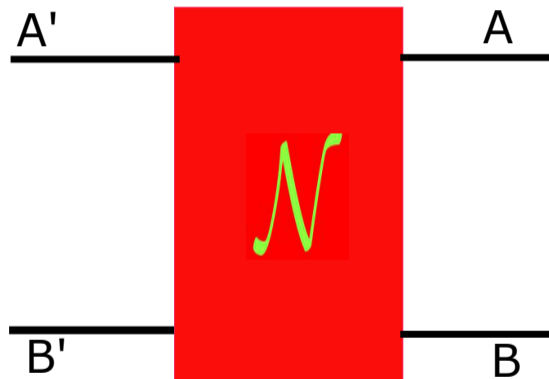


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# Motivation

Bidirectional channels:

- Simple model of quantum network with 2 clients, Alice and Bob.
- Model for quantum gates – CNOT, SWAP, etc.– in noisy intermediate-scale quantum (NISQ) computers.

Entanglement may increase, decrease or not change due to bipartite quantum interactions.

Entanglement distillation: Maximally entangled states are useful resource for several information processing tasks: quantum key distribution, quantum teleportation, etc.

Secret key distillation: Need for secure communication protocols between two parties over network – private reading.

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- Two different information-processing tasks relevant for bipartite quantum interactions:
  - ① Entanglement distillation: generation of singlet state from two separated systems.
  - ② Secret key agreement: generation of maximal classical correlation between two separated systems, such that there's no correlation with the bath.
- New secure communication protocol between two parties, called private reading.

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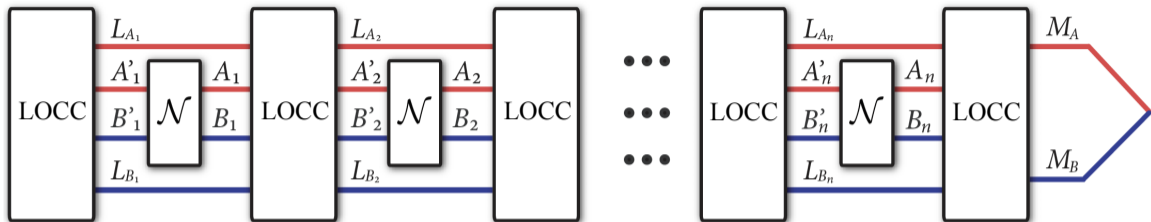
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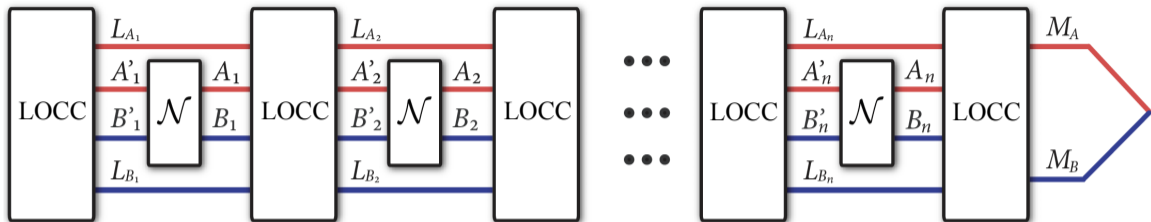
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- Bidirectional max-relative entropy of entanglement:

$$E_{\max}^{2 \rightarrow 2}(\mathcal{N}_{A'B' \rightarrow AB}) = \sup_{\psi_{S_A A'} \otimes \varphi_{B' S_B}} E_{\max}(S_{AA}; S_{B'})_{\mathcal{N}(\psi \otimes \varphi)},$$

where  $\psi_{S_A A'}, \varphi_{B' S_B}$  are pure states,  $S_A \simeq A', S_B \simeq B'$ ,

$E_{\max}(A : B)_\rho = \min_{\sigma_{AB} \in \text{SEP}} D_{\max}(\rho \| \sigma)$ , such that  $D_{\max}(\rho \| \sigma) = \inf \{ \lambda : \rho \leq 2^\lambda \sigma \}$ .

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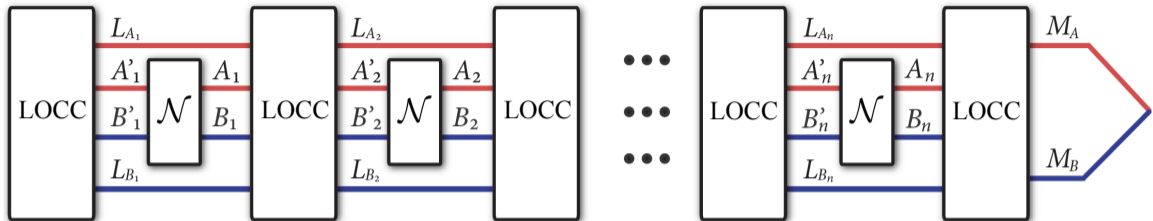
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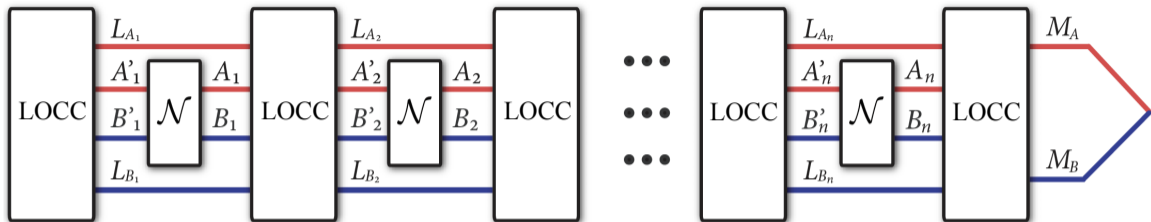
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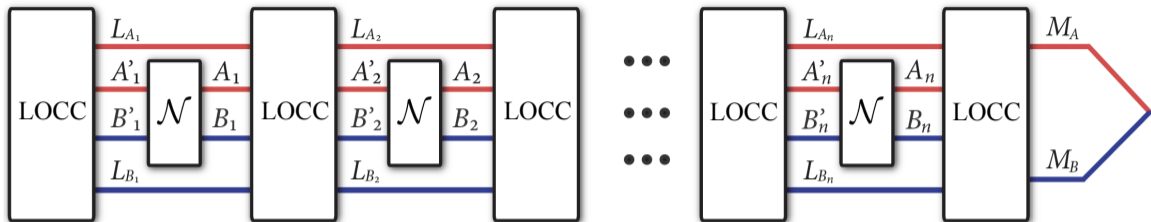


Theorem (LOCC-assisted secret key agreement)

$$\frac{1}{n} \log_2 M \leq E_{\max}^{2 \rightarrow 2}(\mathcal{N}) + \frac{1}{n} \log_2 \left( \frac{1}{1 - \varepsilon} \right).$$

- $P_{LOCC}^{2 \rightarrow 2}(\mathcal{N}_{A'B' \rightarrow AB}) \leq E_{\max}^{2 \rightarrow 2}(\mathcal{N}_{A'B' \rightarrow AB})$ , and this upper bound is in fact a strong converse bound.

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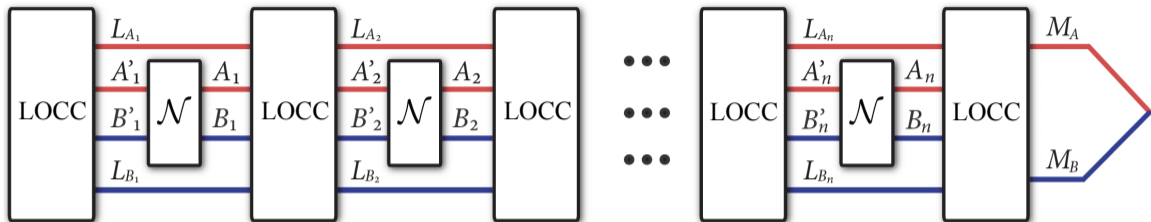


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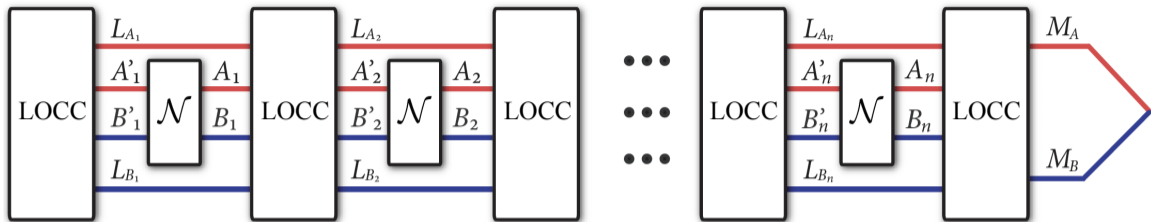
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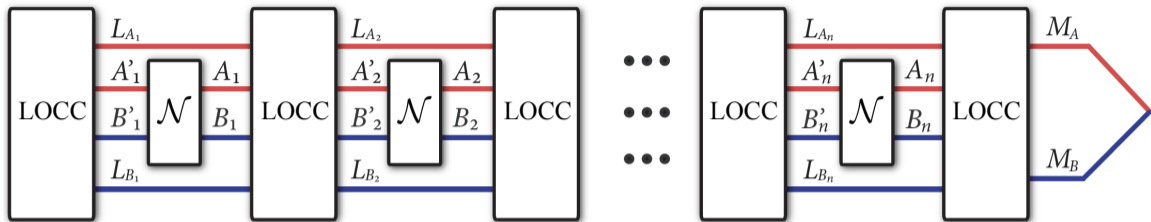


- Private states [HHHO05,HHHO09]: State  $\gamma_{S_A K_A:K_B S_B}$  containing  $\log_2 K$  private bits.
- Success probability in privacy test:  $\text{Tr}\{\Pi\gamma\omega\} \geq 1 - \varepsilon$ . By [HHHO05,HHHO09],  $\text{Tr}\{\Pi\gamma\sigma\} \leq \frac{1}{K}$  for  $\sigma \in \text{SEP}$ .
- Main observation:  $E_{\max}^{2 \rightarrow 2}$  is not enhanced by amortization.

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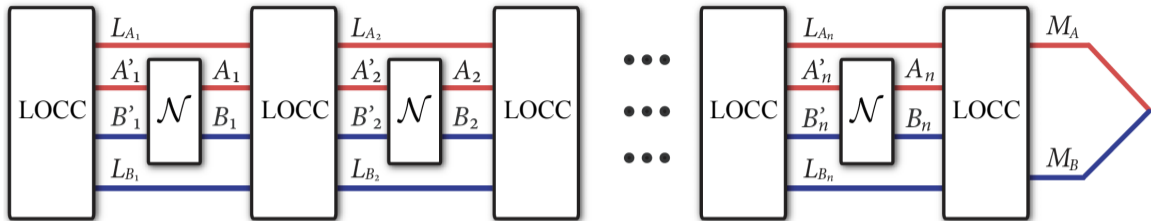


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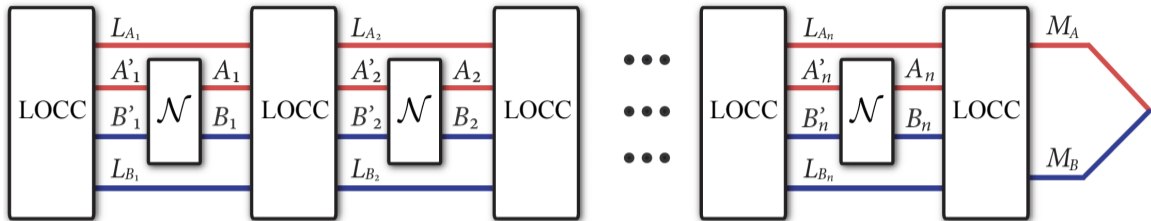


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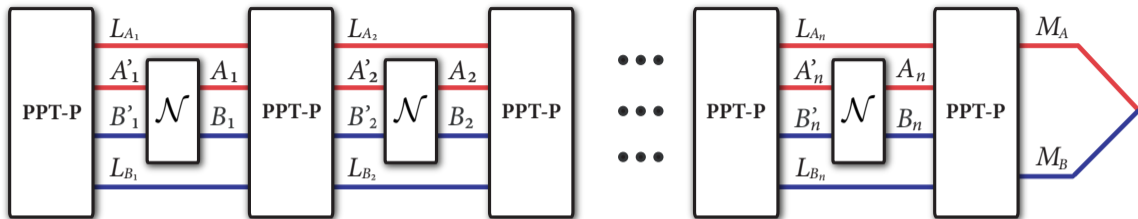


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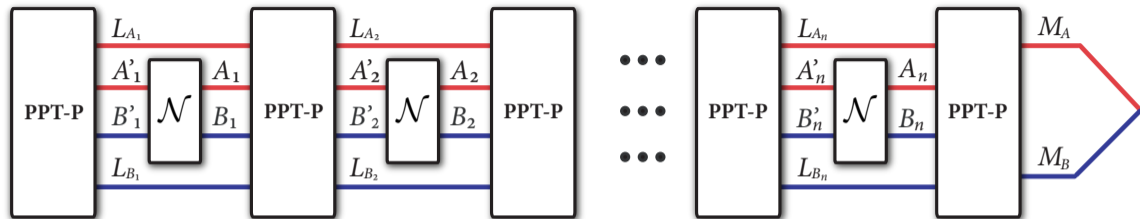
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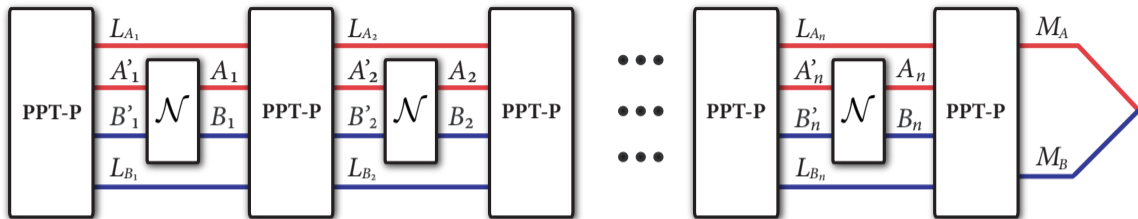
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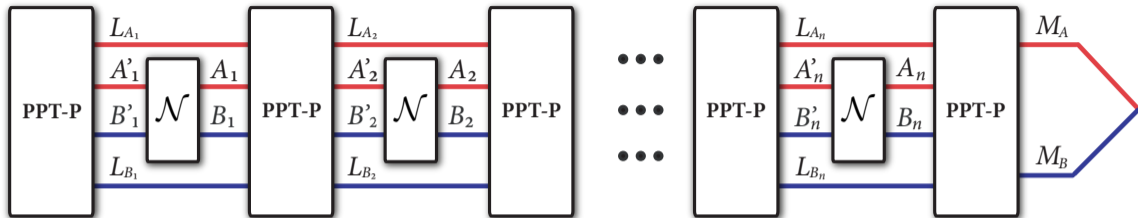
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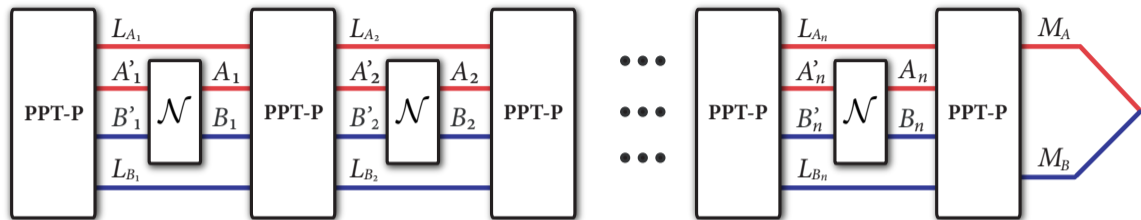
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$$T_{B S_B}\{V_{S_A A B S_B} - Y_{S_A A B S_B}\} \geq J_{S_A A B S_B}^{\mathcal{N}},$$

such that  $S_A \simeq A', S_B \simeq B'$ .



# Entanglement generation over bidirectional channel

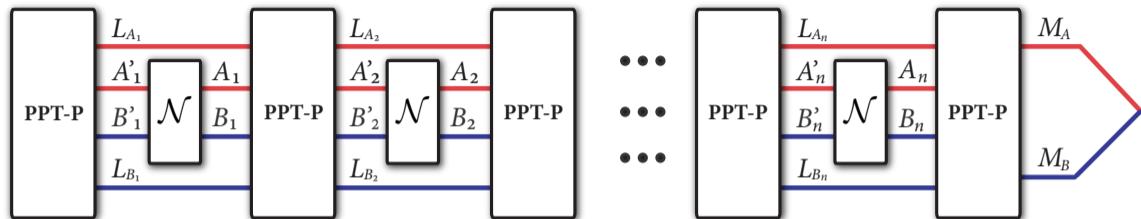


Theorem (PPT-assisted distillable entanglement generation)

$$\frac{1}{n} \log_2 M \leq R_{\max}^{2 \rightarrow 2}(\mathcal{N}) + \frac{1}{n} \log_2 \left( \frac{1}{1 - \varepsilon} \right).$$

- $Q_{PPT}^{2 \rightarrow 2}(\mathcal{N}_{A'B' \rightarrow AB}) \leq R_{\max}^{2 \rightarrow 2}(\mathcal{N}_{A'B' \rightarrow AB})$ , and this upper bound is in fact a strong converse bound.

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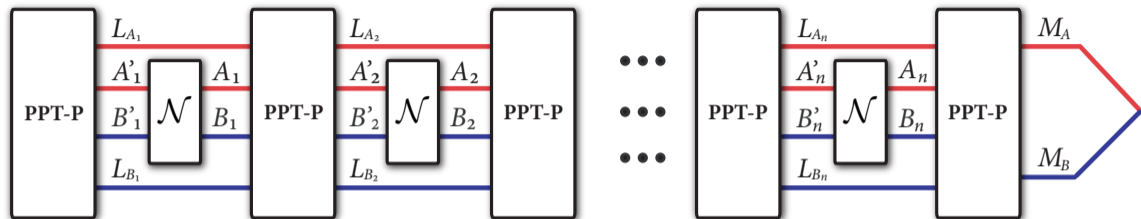


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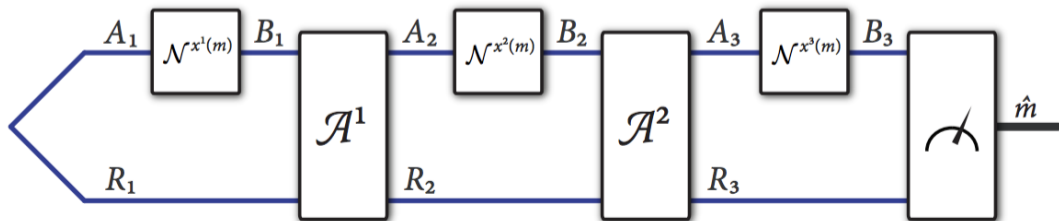
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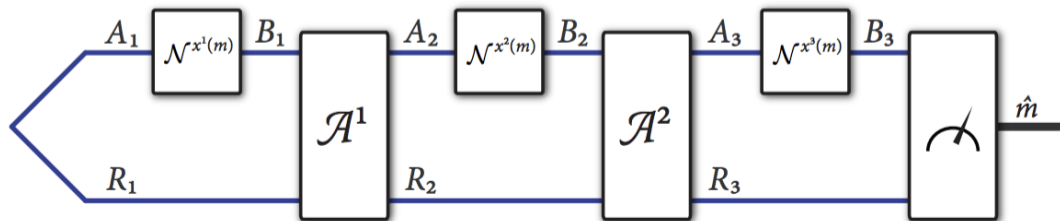
## Application: Private Reading

- First recall: (Quantum) Reading [BRV00,Pir11] of memory devices.
- Memory device: message encoded into a sequence of channels from a memory cell  $\mathcal{S}_{\mathcal{X}} = \{\mathcal{N}_{B' \rightarrow B}^x\}_{x \in \mathcal{X}}$ .
- Alice encodes  $m \in \mathcal{M}$  into codewords  $(x_1(m), \dots, x_n(m))$ , and sets the device to  $(\mathcal{N}_{B' \rightarrow B}^{x_1(m)}, \dots, \mathcal{N}_{B' \rightarrow B}^{x_n(m)})$ .
- Bob can enter quantum states and do channel discrimination to learn  $m$ . Natural to employ adaptive strategy [DW17].



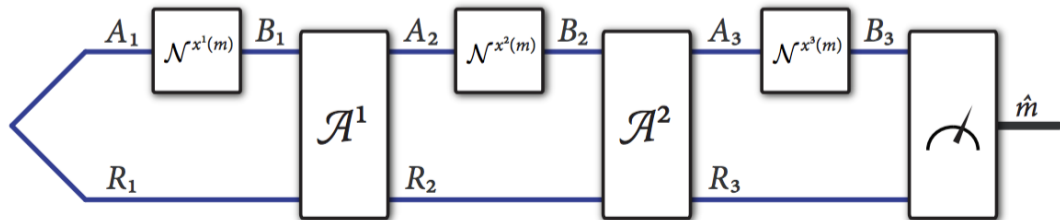
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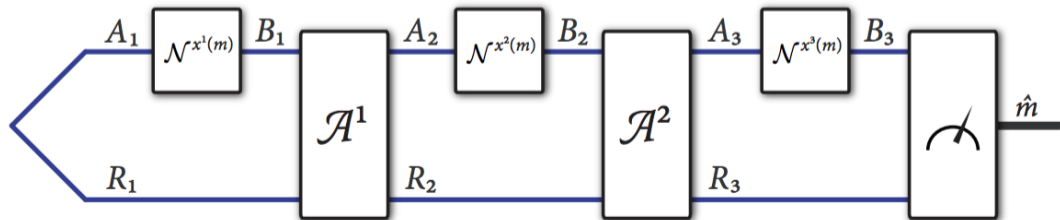
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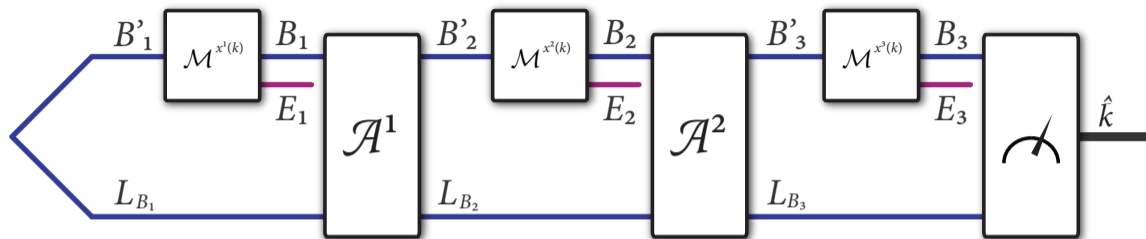


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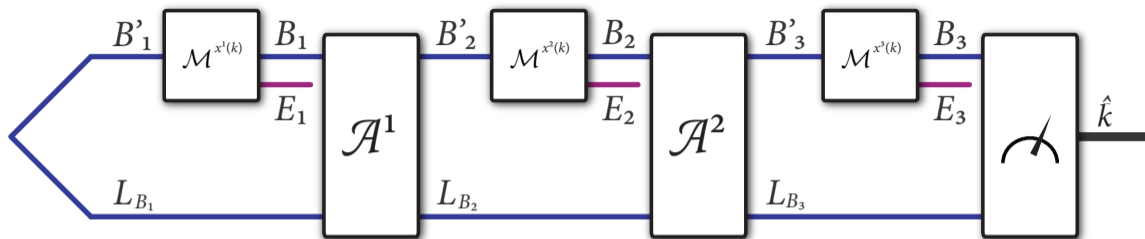
- Private reading: Eve present when Bob performs the readout: Wiretap memory cell  $\bar{\mathcal{S}}_{\mathcal{X}} = \{\mathcal{M}_{B' \rightarrow BE}^x\}_{x \in \mathcal{X}}$ . Special case: Isometric wiretap memory cell  $\bar{\mathcal{S}}_{\mathcal{X}}^{\text{iso}}$ .
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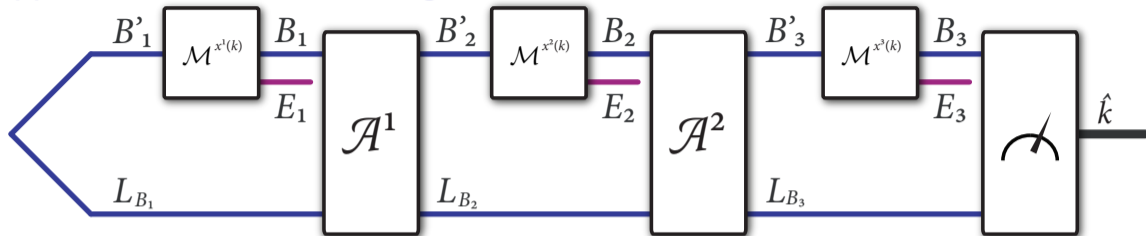


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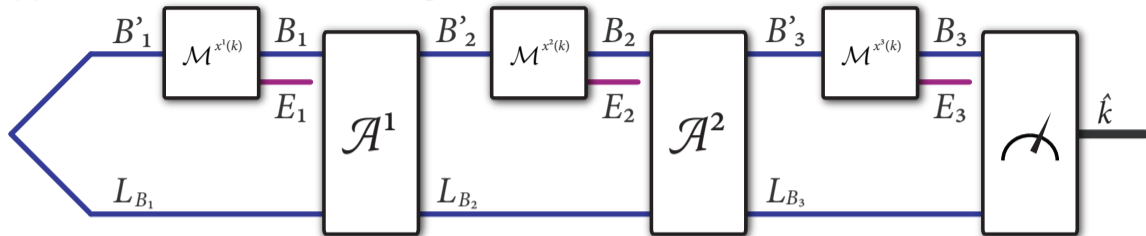


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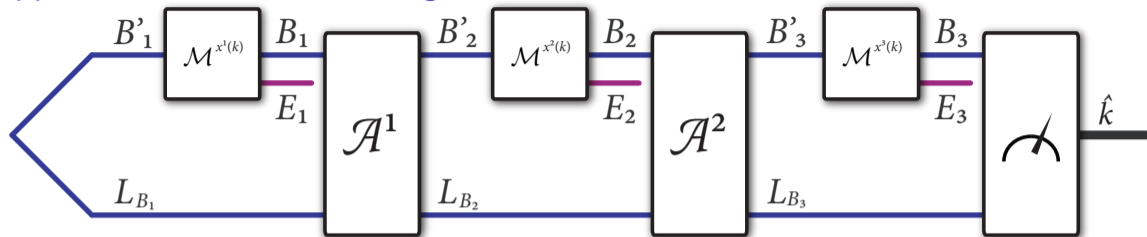


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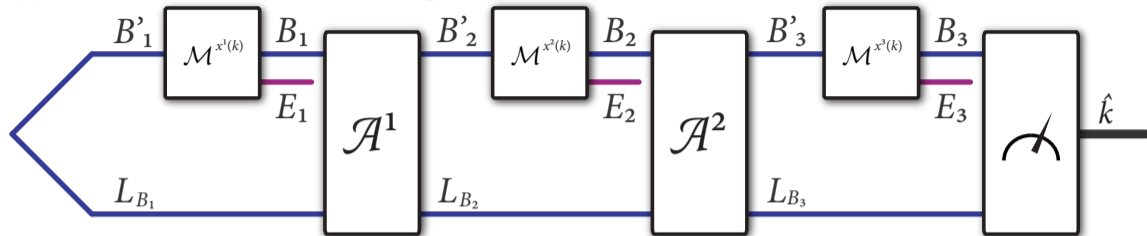


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## Private reading capacity

The strong converse private reading capacity  $\tilde{P}^{\text{read}}(\bar{\mathcal{S}}_{\mathcal{X}}^{\text{iso}})$  of an isometric wiretap memory cell  $\bar{\mathcal{S}}_{\mathcal{X}}^{\text{iso}} = \{U_{B' \rightarrow BE}^{M^x}\}_{x \in \mathcal{X}}$  is bounded from above as

$$\tilde{P}^{\text{read}}(\bar{\mathcal{S}}_{\mathcal{X}}^{\text{iso}}) \leq E_{\max}^{2 \rightarrow 2}(\mathcal{N}_{XB' \rightarrow XB}^{\bar{\mathcal{S}}}),$$

where

$$\mathcal{N}_{XB' \rightarrow XB}^{\bar{\mathcal{S}}}(\cdot) := \text{Tr}_E \left\{ U_{XB' \rightarrow XBE}^{\bar{\mathcal{S}}}(\cdot) \left( U_{XB' \rightarrow XBE}^{\bar{\mathcal{S}}} \right)^\dagger \right\},$$

such that

$$U_{XB' \rightarrow XBE}^{\bar{\mathcal{S}}} := \sum_{x \in \mathcal{X}} |x\rangle\langle x|_X \otimes U_{B' \rightarrow BE}^{M^x}.$$

# Conclusion

- We derived upper bounds on entanglement generation and secret-key-agreement capacities over bidirectional channels. Sizes of reference systems are same as size of input systems (Open question in [BHLS03]).
- Obtain tighter upper bounds for channels obeying certain symmetries, see [DBW17].
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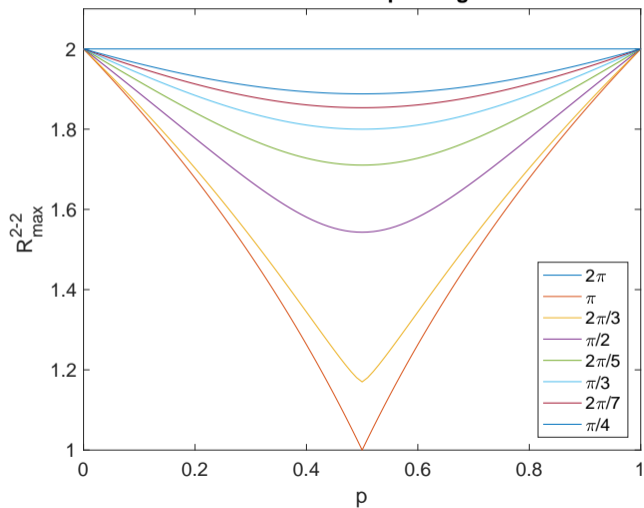
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# SWAP and Collective dephasing

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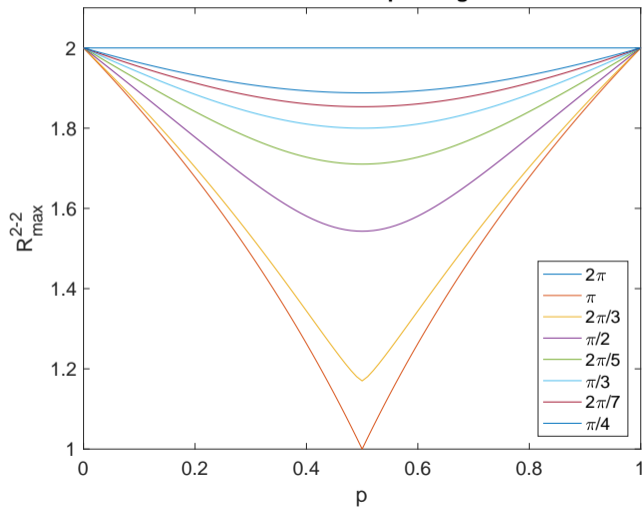


- Collective dephasing:  $|00\rangle \rightarrow |00\rangle$ ,  $|01\rangle \rightarrow e^{i\phi}|01\rangle$ ,  $|10\rangle \rightarrow e^{i\phi}|10\rangle$ ,  $|11\rangle \rightarrow e^{2i\phi}|11\rangle$ .
- Swap operator  $S = \sum_{ij} |ij\rangle\langle ji|$  and collective dephasing:

$$\mathcal{N}_{A'B' \rightarrow AB}(\rho) = pS\rho S^\dagger + (1-p)U^\phi S\rho S^\dagger U^{\phi\dagger}$$

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